

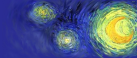
New notion of a renormalized time-ordered product in pAQFT

Katarzyna Rejzner¹

II. Institute for Theoretical Physics, Hamburg University
University of Rome Tor Vergata

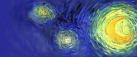
PM 2012, Madrid, 04.06.2012

¹Based on the joint work with Klaus Fredenhagen.



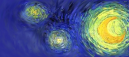
Outline of the talk

- 1 The functional approach
- 2 Free scalar field
- 3 Interaction and renormalization



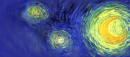
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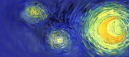
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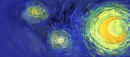
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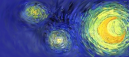
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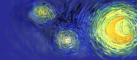
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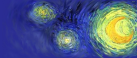
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- A topological $*$ - algebra is a topological vector space, which is a $*$ - algebra and all the algebraic operations are continuous.



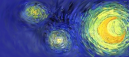
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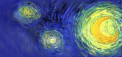
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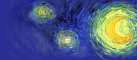


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 - Action:** a map $S_M : \mathcal{D}(M) \rightarrow \mathcal{C}^\infty(\mathcal{E}(M), \mathbb{R})$, where $\mathcal{D}(M) \equiv \mathcal{C}_c^\infty(M, \mathbb{R})$ are compactly supported smooth functions. An example action:

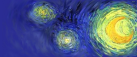
$$S_M(f)(\varphi) = \frac{1}{2} \int (\nabla_\mu \varphi \nabla^\mu \varphi - m^2 \varphi^2)(x) f(x) d\mu(x), \text{ where}$$

$f \in \mathcal{D}(M)$ (cutoff), $\varphi \in \mathcal{E}(M)$, $\mu(x)$ is the volume form on M .



Functionals

- What is $\mathcal{C}^\infty(\mathcal{E}(M), \mathbb{R})$?

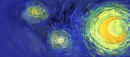


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$$F^{(1)}(\varphi)(h) \doteq \lim_{t \rightarrow 0} \frac{1}{t} (F(\varphi + th) - F(\varphi)) \quad (\text{if exists})$$

F is called **differentiable** if $F^{(1)}(\varphi)(h)$ exists $\forall \varphi \in U, h \in \mathcal{E}(M)$.
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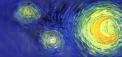
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- The **support** of $F \in \mathcal{C}^\infty(\mathcal{E}(M), \mathbb{R})$ is defined as:

$$\text{supp } F = \{x \in M \mid \forall \text{ neighbourhoods } U \text{ of } x \exists \varphi, \psi \in \mathcal{E}(M), \\ \text{supp } \psi \subset U \text{ such that } F(\varphi + \psi) \neq F(\varphi)\} .$$



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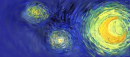
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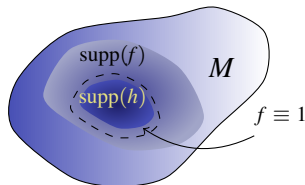
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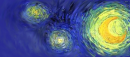
- F is **local** if it is of the form: $F(\varphi) = \int_M f(j_x(\varphi)) d\mu(x)$, where f is a function on the jet bundle over M and $j_x(\varphi) = (\varphi(x), \partial\varphi(x), \dots)$ is the jet of φ at the point x .



Equations of motion

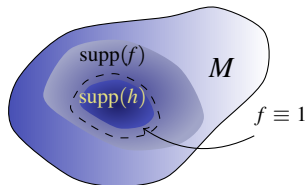
- The Euler-Lagrange derivative of S_M is a map $S'_M : \mathcal{E}(M) \rightarrow \mathcal{D}'(M)$ defined as $\langle S'_M(\varphi), h \rangle = \langle L_M(f)^{(1)}(\varphi), h \rangle$, where $f \equiv 1$ on $\text{supp} h$.



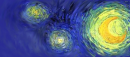


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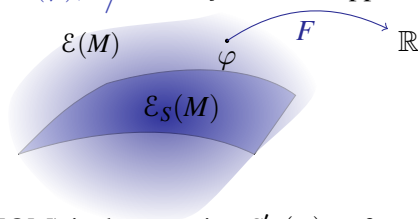


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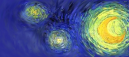


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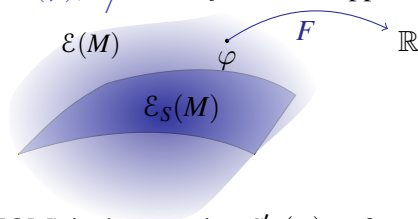


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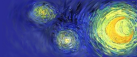


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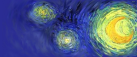


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- We consider the space of smooth functionals on $\mathcal{E}(M)$, i.e. $\mathcal{C}^\infty(\mathcal{E}(M), \mathbb{R})$.



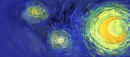
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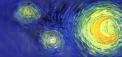
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- Note that $F^{(n)}(\varphi)$ is an element of $\mathcal{E}'(M^n, \mathbb{R})$. We can impose some regularity restrictions on these distributions. In particular we want to control their WF sets.



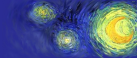
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Definition

Let $u \in \mathcal{D}'(\Omega)$, the wavefront set $\text{WF}(u)$ is the complement in $\Omega \times \mathbb{R}^n \setminus \{0\}$ of the set of $(x, \xi) \in \Omega \times \mathbb{R}^n \setminus \{0\}$ such that there exist $f \in \mathcal{D}(\Omega)$ with $f(x) = 1$ and an open conic neighborhood C of ξ , with

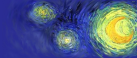
$$\sup_{\xi \in C} (1 + |\xi|)^N |f \cdot \widehat{u}(\xi)| < \infty \quad \forall N \in \mathbb{N}_0.$$



Remarks



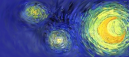
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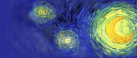
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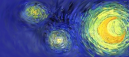


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- In a Lorentzian manifold there are some distinguished directions in TM and T^*M determined by the causal structure (i.e. causal, spacelike, timelike).
- We use these facts to define the algebra $\mathfrak{A}(M)$ by introducing a \star -product on a certain subspace of $\mathcal{C}^\infty(\mathcal{E}(M), \mathbb{R})$.



Free scalar field

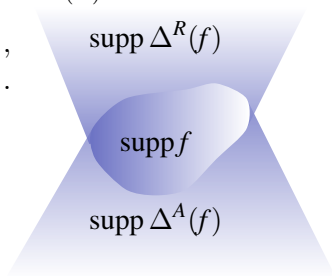
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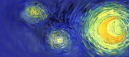


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 $P \circ \Delta^{R/A} = \text{id}_{\mathcal{D}(M)}$, $\Delta^{R/A} \circ (P|_{\mathcal{D}(M)}) = \text{id}_{\mathcal{D}(M)}$ and

$$\text{supp}(\Delta^R) \subset \{(x, y) \in M^2 | y \in (\overline{V}_-)_x\},$$
$$\text{supp}(\Delta^A) \subset \{(x, y) \in M^2 | y \in (\overline{V}_+)_x\}.$$



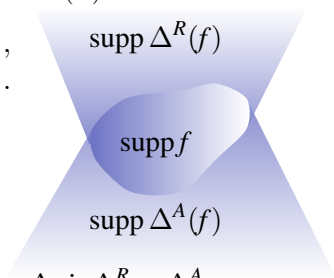


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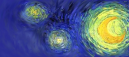
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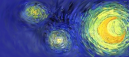


- Their difference is the causal propagator $\Delta \doteq \Delta^R - \Delta^A$.



Propagators and Green's functions

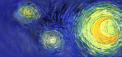
- $\text{WF}(\Delta) = \{(x, k; x', -k') \in \dot{T}^*M^2 \mid (x, k) \sim (x', k')\}$, where \sim means that there is a causal curve connecting x and x' and k' is the parallel transport of k along it.



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- We can always decompose Δ to positive and negative frequency parts: $i\Delta = \Delta_+ + \Delta_-$, i.e.:

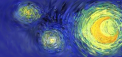
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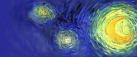
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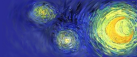
- We can also define the Feynman propagator:

$$\Delta_F = \frac{i}{2}(\Delta^A + \Delta^R) + \Delta_1.$$



★-product

- Properties of the WF set of Δ_+ motivate the following definition:



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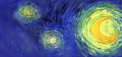
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$$\Xi_n \doteq T^*M^n \setminus \{(x_1, \dots, x_n, k_1, \dots, k_n) \mid k_i \in (\bar{V}_+)_{x_i} \cup (\bar{V}_-)_{x_i}, i = 1 \dots n\}.$$



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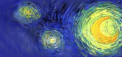
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- We define the ★-product (deformation of the pointwise product):

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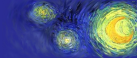
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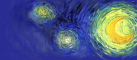
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- The free QFT is defined as $\mathfrak{A}_0(M) \doteq (\mathfrak{F}_{\text{mc}}(M), \star, *)$, where $F^*(\varphi) \doteq \overline{F(\varphi)}$.



Time-ordered product

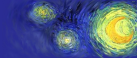
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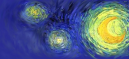
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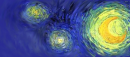
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- Define the time-ordered product $\cdot_{\mathcal{T}}$ on $\mathcal{T}(\mathfrak{F}_{\text{reg}}(M)[[\hbar]])$ by:

$$F \cdot_{\mathcal{T}} G \doteq \mathcal{T}(\mathcal{T}^{-1}F \cdot \mathcal{T}^{-1}G)$$

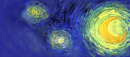


Interaction

- We now have an algebraic structure with two products $(\mathfrak{F}_{\text{reg}}(M)[[\hbar]], \star, \cdot_{\mathcal{T}})$, where \star is non-commutative, $\cdot_{\mathcal{T}}$ is commutative and they are related by a causal relation:

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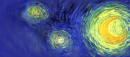
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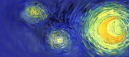
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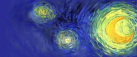
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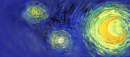
- Because of the WF set properties of Δ_F , the time-ordered product $\cdot_{\mathcal{T}}$ is not well defined on local, non-constant functionals, but the physical interaction is usually local!



Renormalization problem

- **Renormalization problem:** extend $\mathcal{S}(\cdot)$ to local arguments. This is reduced to extending the n -fold time-ordered products, since we can define:

$$\mathcal{S}(V) = \sum_{n=0}^{\infty} \frac{1}{n!} \mathcal{T}^n(V, \dots, V).$$

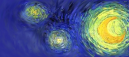


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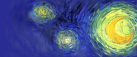


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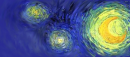
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- To extend \mathcal{T}^n to arbitrary local functionals we use the causal approach of Epstein and Glaser (causal perturbation theory).



Causal perturbation theory

In causal perturbation theory n -fold time ordered products have to obey following axioms:

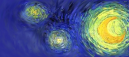
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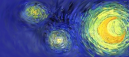


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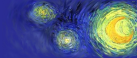
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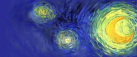
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By the theorem of Epstein and Glaser we know that the **extension exists, but is not unique**. The theorem is proved inductively (in n) and at each step the problem is reduced to an extension of a real valued distribution.



Causal perturbation theory

- We can split the time-ordered product, \mathcal{T}^n , into two parts:

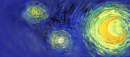


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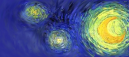
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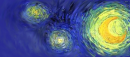
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- The n -fold time-ordered product can then be written as

$$F_1 \cdot_{\mathcal{T}} \cdots \cdot_{\mathcal{T}} F_n = \sum_{\alpha \in \mathbb{N}^n} \left\langle \widetilde{S}_\alpha, \delta^\alpha (F_1 \otimes \cdots \otimes F_n) \right\rangle$$

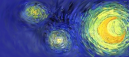
where we denoted:

$$\widetilde{S}_\alpha \doteq \sum_{\Gamma \in \mathcal{G}_\alpha} \frac{\hbar^{|\Gamma|}}{\text{Sym}(\Gamma)} \widetilde{S}_\Gamma.$$



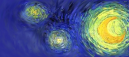
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- The distribution \widetilde{S}_α is a sum over \mathcal{G}_α , the set of (non-tadpole) graphs with $n = \dim(\alpha)$ vertices and $\frac{|\alpha|}{2}$ lines such that there are α_i lines joining at vertex i and $\text{Sym}(\Gamma) \in \mathbb{N}$ is the so called symmetry factor of the graph Γ .



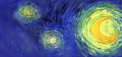
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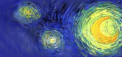
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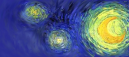
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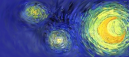
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$$S(V) = \sum_{n=0}^{\infty} \frac{1}{n!} \mathcal{T}^n(V, \dots, V)$$
- Let $Z \in \mathcal{R}$, then we obtain a new S-matrix as $\widehat{S} = S \circ Z$ and all the S-matrices are obtained in this way.



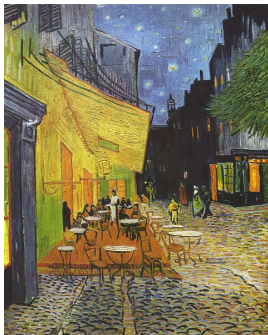
Renormalized n -fold time-ordered products



- We have already maps
$$\mathcal{T}^n : \mathfrak{F}_{\text{loc}}(M)^{\otimes n} \rightarrow \mathfrak{F}_{\text{mc}}(M)[[\hbar]],$$
but we can get even more! There exists a map
$$\beta : \mathfrak{F}(M) \rightarrow S^\bullet \mathfrak{F}_{\text{loc}}^{(0)}(M),$$
inverse to the pointwise multiplication.



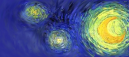
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$$F \cdot_{\mathcal{T}_r} G \doteq \mathcal{T}_r(\mathcal{T}_r^{-1}F \cdot \mathcal{T}_r^{-1}G),$$

which is an associative, commutative product on $\mathcal{T}_r(\mathfrak{F}(M)) \subset \mathfrak{A}_0(M)$.



Renormalized n -fold time-ordered products

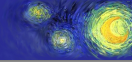







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- **The renormalized QFT** is a structure with two products $\mathfrak{A}(M) = (\mathcal{T}_r(\mathfrak{F}(M)), \star, \cdot_{\mathcal{T}_r})$.



-  R. Brunetti, K. Fredenhagen, *Microlocal analysis and interacting quantum field theories: Renormalization on physical backgrounds*, *Commun. Math. Phys.* **208**, 623 (2000).
-  K. J. Keller, *Dimensional Regularization in Position Space and a Forest Formula for Regularized Epstein-Glaser Renormalization*, Ph.D thesis, Hamburg 2010, [arXiv:math-ph/1006.2148v1].
-  Ch. Bär, K. Fredenhagen, *Quantum field theory on curved spacetimes*, Lecture Notes in Physics, Vol. 786, Springer Berlin Heidelberg 2009.
-  K. Rejzner, *Quantum field theory for mathematicians*, course given 6-11 February 2012 in Hamburg, http://rejzner.com/files/QFT_mat.pdf.
-  K. Fredenhagen, K. Rejzner, *Batalin-Vilkovisky formalism in perturbative algebraic quantum field theory*, [arXiv:math-ph/1110.5232].



Thank you for your attention!